

# Leverrier-Takeno and Faddeev-Sominsky algorithms

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**Abstract** We exhibit representations of the coefficients of the characteristic polynomial of any matrix  $A_{n \times n}$ , especially in terms of the (exponential) complete Bell polynomials. Besides, we use the Faddeev-Sominsky method to obtain the Lanczos formula for the resolvent of a matrix. We indicate that the Newton's recurrence formula can be solved via the inversion of a triangular matrix.

**Keywords** Faddeev-Sominsky's algorithm, Cayley-Hamilton-Frobenius theorem, Leverrier-Takeno's procedure, (Exponential) complete Bell polynomials.

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## §1. Introduction

For an arbitrary matrix  $A_{n \times n} = (A_j^i)$  its characteristic polynomial [1-3]:

$$p(\lambda) \equiv \lambda^n + a_1 \lambda^{n-1} + \dots + a_{n-1} \lambda + a_n, \quad (1)$$

can be obtained, through several procedures [1,4-8], directly from the condition  $\det (A_j^i - \lambda \delta_j^i) = 0$ . The method of Leverrier-Takeno [4,9-13] is a simple and interesting technique to construct (1) based in the traces of the powers  $A^r$ ,  $r = 1, \dots, n$ .

On the other hand, it is well known that an arbitrary matrix satisfies its characteristic equation, which is the Cayley-Hamilton-Frobenius identity [1-3,14]:

$$A^n + a_1 A^{n-1} + \dots + a_{n-1} A + a_n I = 0. \quad (2)$$

If  $A$  is non-singular (that is,  $\det A \neq 0$ ), then from (2) we obtain its inverse matrix:

$$A^{-1} = -\frac{1}{a_n} \left( A^{n-1} + a_1 A^{n-2} + \dots + a_{n-1} I \right), \quad (3)$$

where  $a_n \neq 0$  because  $a_n = (-1)^n \det A$ .

Faddeev-Sominsky [15–24] proposed an algorithm to determine  $A^{-1}$  in terms of  $A^r$  and their traces, which is equivalent [23] to the Cayley-Hamilton-Frobenius theorem (2) plus the Leverrier-Takeno's method to construct the characteristic polynomial of a matrix  $A$ . In Sec. 2, we use the Faddeev-Sominsky's procedure to deduce the Lanczos expression [25] for the resolvent of  $A$  [20, 21, 26, 27], that is, the Laplace transform of  $\exp(tA)$  [28].

## §2. Leverrier-Takeno technique

If we define the quantities:

$$a_0 = 1, \quad s_k = \text{tr } A^k \quad k = 1, 2, \dots, n \quad (4)$$

then the approach of Leverrier-Takeno [4, 9–13] implies (1) wherein the  $a_i$  are determined with the Newton's recurrence relation:

$$ra_r + s_1 a_{r-1} + s_2 a_{r-2} + \dots + s_{r-1} a_1 + s_r = 0, \quad r = 1, 2, \dots, n \quad (5)$$

therefore:

$$\begin{aligned} a_1 &= -s_1, & 2! a_2 &= (s_1)^2 - s_2, & 3! a_3 &= -(s_1)^3 + 3s_1 s_2 - 2s_3, \\ 4! a_4 &= (s_1)^4 - 6(s_1)^2 s_2 + 8s_1 s_3 + 3(s_2)^2 - 6s_4, \end{aligned} \quad (6)$$

$$5! a_5 = -24s_5 - (s_1)^5 + 10(s_1)^3 s_2 - 20(s_1)^2 s_3 - 15(s_2)^2 s_1 + 30s_1 s_4 + 20s_2 s_3, \text{ etc,}$$

in particular,  $\det A = (-1)^n a_n$ , that is, the determinant of any matrix only depends on the traces  $s_r$ , which means that  $A$  and its transpose have the same determinant. In [29–31] we find the general expression:

$$a_k = \frac{(-1)^k}{k!} \begin{vmatrix} s_1 & k-1 & 0 & \dots & 0 \\ s_2 & s_1 & k-2 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots \\ s_{k-1} & s_{k-2} & \dots & \dots & 1 \\ s_k & s_{k-1} & \dots & \dots & s_1 \end{vmatrix}, \quad k = 1, \dots, n, \quad (7)$$

which allows reproduce the values (6).

We can exhibit a relation to determine the coefficients  $a_j$  via the complete Bell polynomials [8, 32–40], in fact, we have the following representation [8]:

$$m! a_m = Y_m \left( -0!s_1, -1!s_2, -2!s_3, -3!s_4, \dots, -(m-2)!s_{m-1}, -(m-1)!s_m \right). \quad (8)$$

such that [37, 41]:

$$Y_m(x_1, x_2, \dots, x_m) = \begin{pmatrix} \binom{m-1}{0}x_1 & \binom{m-1}{1}x_2 & \dots & \binom{m-1}{m-2}x_{m-1} & \binom{m-1}{m-1}x_m \\ -1 & \binom{m-2}{0}x_1 & \dots & \binom{m-2}{m-3}x_{m-2} & \binom{m-2}{m-2}x_{m-1} \\ 0 & -1 & \dots & \binom{m-3}{m-4}x_{m-3} & \binom{m-3}{m-3}x_{m-2} \\ \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & \dots & \binom{1}{0}x_1 & \binom{1}{1}x_2 \\ 0 & 0 & \dots & -1 & \binom{0}{0}x_1 \end{pmatrix}, \quad (9)$$

therefore:

$$\begin{aligned} Y_0 &= 1, \quad Y_1 = x_1, \quad Y_2 = x_1^2 + x_2, \\ Y_3 &= x_1^3 + 3 x_1 x_2 + x_3, \quad Y_4 = x_1^4 + 6 x_1^2 x_2 + 4 x_1 x_3 + 3 x_2^2 + x_4, \\ Y_5 &= x_1^5 + 10 x_1^3 x_2 + 10 x_1^2 x_3 + 15 x_1 x_2^2 + 5 x_1 x_4 + 10 x_2 x_3 + x_5, \dots \end{aligned} \quad (10)$$

We see that (8) and (10) imply (6) if we employ  $x_1 = -s_1, x_2 = -s_2, x_3 = -2s_3, x_4 = -6s_4, x_5 = -24s_5, \dots$ . It is simple to prove that (9) with  $x_k = -(k-1)! s_k$  gives (7), thus the coefficients of the characteristic polynomial (1) are generated by the (exponential) complete Bell polynomials.

In the Newtons formula (5) the quantities  $s_r$  are known, and the  $a_j$  are solutions of the triangular linear system [42–44]:

$$A_{n \times n} (a_j)_{n \times 1} \equiv \begin{pmatrix} 1 & 0 & 0 & \dots & \dots & 0 \\ s_1 & 2 & 0 & \dots & \dots & 0 \\ s_2 & s_1 & 3 & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots & 0 \\ s_{n-1} & s_{n-2} & s_{n-3} & \dots & s_1 & n \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \dots \\ \dots \\ \dots \\ a_n \end{pmatrix} = - \begin{pmatrix} s_1 \\ s_2 \\ s_3 \\ \dots \\ \dots \\ \dots \\ s_n \end{pmatrix}, \quad (11)$$

then:

$$\begin{pmatrix} a_1 \\ \dots \\ \dots \\ a_n \end{pmatrix} = -A^{-1} \begin{pmatrix} s_1 \\ \dots \\ \dots \\ s_n \end{pmatrix}, \quad (12)$$

which gives the opportunity to invert a triangular matrix via interesting algorithms applying the Faddeev-Sominsky method [15–24], matrix multiplication [45, 46] or binomial series [47].

### §3. Faddeev-Sominsky procedure

The Faddeev-Sominsky's technique to obtain  $A^{-1}$  is a sequence of algebraic computations on the powers  $A^r$  and their traces, in fact, this algorithm is given via the following instructions:

$$\begin{aligned}
 B_1 &= A, & q_1 &= \operatorname{tr} B_1, & C_1 &= B_1 - q_1 I, \\
 B_2 &= C_1 A, & q_2 &= \frac{1}{2} \operatorname{tr} B_2, & C_2 &= B_2 - q_2 I, \\
 &\dots & &\dots & &\dots \\
 &\dots & &\dots & &\dots \\
 B_{n-1} &= C_{n-2} A, & q_{n-1} &= \frac{1}{n-1} \operatorname{tr} B_{n-1}, & C_{n-1} &= B_{n-1} - q_{n-1} I, \\
 B_n &= C_{n-1} A, & q_n &= \frac{1}{n} \operatorname{tr} B_n,
 \end{aligned} \tag{13}$$

then:

$$A^{-1} = \frac{1}{q_n} C_{n-1}. \tag{14}$$

For example, if we apply (13) for  $n = 4$ , then it is easy to see that the corresponding  $q_r$  imply (6) with  $q_j = -a_j$ , and besides (14) reproduces (3). By mathematical induction one can prove that (13) and (14) are equivalent to (3), (4) and (5), showing [23] thus that the Faddeev-Sominsky's technique has its origin in the Leverrier-Takeno method plus the Cayley-Hamilton-Frobenius theorem.

From (13) we can see that [26]:

$$C_k = A^k + a_1 A^{k-1} + a_2 A^{k-2} + \dots + a_{k-1} A + a_k I, \quad k = 1, 2, \dots, n-1, \quad C_n = B_n - q_n I = 0, \tag{15}$$

and for  $k = n - 1$ :

$$C_{n-1} = A^{n-1} + a_1 A^{n-2} + a_2 A^{n-3} + \dots + a_{n-2} A + a_{n-1} I \stackrel{(3)}{=} -a_n A^{-1}$$

in harmony with (14) because  $a_n = -q_n$ . The property  $C_n = 0$  is equivalent to (2); if  $A$  is singular, the process (13) gives the adjoint matrix of  $A$  [2, 3, 16], in fact,  $\operatorname{Adj} A = (-1)^{n+1} C_{n-1}$ .

If the roots of (1) have distinct values, then the Faddeev-Sominsky's algorithm allows obtain the corresponding eigenvectors of  $A$  [6]:

$$A \vec{u}_k = \lambda_k \vec{u}_k, \quad k = 1, 2, \dots, n, \tag{16}$$

because for a given value of  $k$ , each column of:

$$Q_k \equiv \lambda_k^{n-1} I + \lambda_k^{n-2} C_1 + \dots + C_{n-1}. \tag{17}$$

satisfies (16) [16, 18, 27], and therefore all columns of  $Q_k$  are proportional to each other, that is,  $\text{rank } Q_k = 1$  [18]; we note that  $Q_k = Q(\lambda_k)$  with the participation of the matrix:

$$Q(z) \equiv z^{n-1}I + z^{n-2}C_1 + z^{n-3}C_2 + \dots + zC_{n-2} + C_{n-1}. \quad (18)$$

By synthetic division of two polynomials [1]:

$$\frac{p(z)}{z - \lambda} = \sum_{r=0}^{n-1} \left( \lambda^r + a_1\lambda^{r-1} + a_2\lambda^{r-2} + \dots + a_{r-1}\lambda + a_r \right) z^{n-1-r},$$

then under the change  $\lambda \rightarrow A$  we obtain the Lanczos expression [25] for the resolvent of a matrix [20, 21, 26, 27]:

$$\frac{1}{zI - A} = \frac{1}{p(z)} \sum_{r=0}^{n-1} z^{n-1-r} C_r = \frac{Q(A)}{p(z)}, \quad (19)$$

if  $A$  is non-singular, then [19] for  $z = 0$  implies (14). McCarthy [48] used (19) and the Cauchy's integral theorem in complex variable to show the Cayley-Hamilton-Frobenius identity indicated in (2); the relation (19) is the Laplace transform of  $\exp(tA)$  [28].

On the other hand, Sylvester [49–52] obtained the following interpolating definition of  $f(A)$ :

$$f(A) = \sum_{j=1}^n f(\lambda_j) \prod_{k \neq j} \frac{A - \lambda_k I}{\lambda_j - \lambda_k}, \quad (20)$$

which is valid if all eigenvalues are different from each other. Buchheim [53] generalized (20) to multiple proper values using Hermite interpolation, thereby giving the first completely general definition of a matrix function. From (19) and (20) for  $f(s) = 1/(z-s)$  we deduce the properties:

$$Q(z) = \sum_{j=1}^n \prod_{k=1, k \neq j}^n \frac{z - \lambda_k}{\lambda_j - \lambda_k} (A - \lambda_k I), \quad Q_j = \prod_{k=1, k \neq j}^n (A - \lambda_k I)$$

$$Q_j \vec{u}_j = \prod_{k=1, k \neq j}^n (\lambda_j - \lambda_k) \vec{u}_j, \quad (21)$$

hence the eigenvectors of  $A$  showed in (16) also are proper vectors of the matrices  $Q_j$ . Besides, from (16) and (21):

$$A Q_j \vec{u}_j = \prod_{k=1, k \neq j}^n (\lambda_j - \lambda_k) \lambda_j \vec{u}_j = \lambda_j Q_j \vec{u}_j, \quad A Q_j = \lambda_j Q_j \quad (22)$$

that is, each column of  $Q_j$  is eigenvector of  $A$  with proper value  $\lambda_j$ . The resolvent (19) implies the relation  $(A - zI)Q(z) = -p(z)I$ , then  $(A - \lambda_k I) Q(\lambda_k) = -p(\lambda_k)I = 0$  in according with (22).

From the Sylvester's formula (20) with  $f(z) = p(z)$  we obtain  $p(A) = 0$ , which is the Cayley-Hamilton-Frobenius theorem indicated in (2). If  $f(z) = e^{tz}$ , then (20) allows to construct  $\exp(tA)$  that, in particular, is valuable to determine the motion of classical charged particles into a homogeneous electromagnetic field, and to integrate the Frenet-Serret equations with constant curvatures [54]. In [51, 55] we find that the book of Frazer-Duncan-Collar [56]

emphasizes the important role of the matrix exponential in solving differential equations and was the first to employ matrices as an engineering tool, and indeed the first book to treat matrices as a branch of applied mathematics.

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